

Constrained fitting of disentangled binary star spectra: application to V615 Persei in the open cluster h Persei (Research Note)

E. Tamajo¹, K. Pavlovski^{1,2}, and J. Southworth²

¹ Department of Physics, University of Zagreb, Bijenička 32, 10 000 Zagreb, Croatia
e-mail: pavlovski@phy.hr

² Astrophysics Group, Keele University, Newcastle-under-Lyme, ST5 5BG, UK

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ABSTRACT

Context. Using the technique of spectral disentangling, it is possible to determine the individual spectra of the components of a multiple star system from composite spectra observed at a range of orbital phases. This method has several advantages: it is unaffected by line blending, does not use template spectra, and returns individual component spectra with very high signal-to-noise ratios.

Aims. The disentangled spectra of a binary star system are very well suited to spectroscopic analysis but for one problem: the absolute spectral line depths are unknown because this information is not contained in the original spectra (unless there is one taken in eclipse) without making assumptions about the spectral characteristics of the component stars. Here we present a method for obtaining the atmospheric parameters of the component stars by the constrained fitting of synthetic spectra to observed and disentangled spectra.

Methods. Disentangled spectra are fitted using synthetic spectra and a genetic algorithm in order to determine the effective temperatures, surface gravities and relative light contributions of the two stars in a binary system. The method is demonstrated on synthetic spectra and then applied to the eclipsing binary V615 Per, a member of the young open cluster NGC 869 (h Persei).

Results. The method works well for disentangled spectra with signal-to-noise ratios of 100 or more. For V615 Per we find a normal He abundance but an Mg abundance, which indicates bulk metallicity, a factor of two lower than typical for nearby OB stars.

Key words. binaries: spectroscopic – stars: abundances – stars: atmospheres – open clusters and associations: individual: h Persei

1. Introduction

The technique of *spectral disentangling* (SPD) allows the isolation of the individual spectra of the component stars of a double-lined spectroscopic binary system from a set of composite spectra observed over a range of orbital phases. It was originally formulated in the wavelength domain by [Simon & Sturm \(1994\)](#) and in the Fourier domain by [Hadrava \(1995\)](#). The technique simultaneously returns the best-fitting individual spectra and the orbital velocity amplitudes of the two stars. A detailed overview of SPD can be found in [Pavlovski & Hensberge \(2009\)](#).

Compared to other methods of radial velocity measurement, SPD has several advantages. Firstly, it is independent of template spectra so avoids any systematic errors due to spectral differences between the target and template stars. Secondly, it is not affected by the blending of spectral lines of the two stars (see [Hensberge et al. 2000](#); [Southworth & Clausen 2007](#)). Thirdly, the resulting disentangled spectra contain the combined signal of the input spectra ([Pavlovski & Southworth 2009](#)) so have a much higher signal-to-noise (S/N) ratio.

There are two disadvantages of the SPD approach. The first of these is that the continuum normalisation of the input spectra has to be very good in order to avoid low-frequency spurious patterns in the resulting disentangled spectra ([Hensberge et al. 2008](#)). The second is that relative continuum light contributions of the two stars cannot be found using SPD as this information is itself not contained in the observed spectra, unless a spectrum has been obtained during an eclipse ([Ilijić et al. 2004](#)).

SPD is well suited to the spectral analysis of stars in binary systems. Each disentangled spectrum contains only features due to one star, so can be analysed using standard methods for single stars. The high S/N ratios of disentangled spectra are very helpful to this process, but the undetermined continuum light ratio between the component stars complicates the spectral analysis. In this work we present a method to fit synthetic spectra to disentangled spectra, where the atmospheric parameters of the stars are determined simultaneously with the relative light contributions of the stars. A genetic algorithm is used for the optimization in order to ensure that the best solution is found in a parameter space which suffers from strong degeneracies, in particular between effective temperature (T_{eff}) and surface gravity ($\log g$).

An important application of SPD is the study of detached eclipsing binary star systems (dEBs). These represent the primary source of directly-measured masses and radii of stars, and as such are cornerstones of stellar physics ([Andersen 1991](#); [Torres et al. 2010](#)). SPD can be used to measure the velocity amplitudes of the stars, which are necessary for the mass and radius measurements, simultaneously with obtaining the individual stellar spectra for spectral analysis ([Pavlovski & Hensberge 2005](#); [Pavlovski et al. 2009](#); [Pavlovski & Southworth 2009](#)). A major advantage of dEBs to this process is that the surface gravities of the stars can be obtained to within ± 0.01 dex from the mass and radius measurements: these parameters can then be fixed in the spectral analysis and thus the degeneracy between T_{eff} and $\log g$ avoided ([Simon et al. 1994](#); [Hensberge et al. 2000](#)).

In this work we demonstrate the genetic algorithm approach to fitting disentangled spectra on the dEB V615 Persei, a member of the young open cluster h Persei. Southworth et al. (2004a, hereafter SMS04) obtained a series of high-resolution spectra of V615 Per and analysed them with published light curves (Krzesiński et al. 1999) to measure the masses (4.08 and 3.18 $M_{\odot}$) to accuracies of 2% and the radii (2.29 and 1.90 $R_{\odot}$) to 5%, resulting in surface gravities measured to within 0.05 dex. The T_{eff} values were found to be $15\,000 \pm 500$ K and $11\,000 \pm 500$ K. SMS04 found that stellar evolutionary models required a subsolar metal abundance ($Z \approx 0.01$) to reproduce the measured masses and radii of V615 Per.

The Perseus Double Cluster comprises h Persei (NGC 869) and χ Persei (NGC 884). It has been extensively studied via deep CCD photometry (Keller et al. 2001; Marco & Bernabeu 2001; Slesnick et al. 2002; Capilla & Fabregat 2002; Currie et al. 2010), from which there is general agreement on its distance (2.3 to 2.4 kpc) and age (13–14 Myr). But these studies assumed a solar chemical composition, and their results may be systematically wrong if this assumption is incorrect.

Conflicting results on the chemical composition of the Perseus Double Cluster are present in the literature. Detailed abundance analyses based on high-resolution spectra of hot stars (Lennon et al. 1988; Dufton et al. 1990) have challenged previous findings of low helium abundances (Nissen 1976; Klochkova & Panchuk 1987; Wolff & Heasley 1985). Dufton et al. (1990) and Smartt & Rolleston (1997) found an approximately solar metal abundance from high-resolution spectra, but this was not supported by Vrancken et al. (2000). In this work we attempt to shed additional light on this subject by measuring the helium and metal abundances of the stars in the dEB V615 Per.

2. Constrained fitting of disentangled spectra using a genetic algorithm

A computer code has been constructed which fits synthetic spectra to disentangled spectra of a binary system in order to determine the atmospheric parameters T_{eff} , $\log g$, projected rotational velocities $v \sin i$, Doppler shifts, and the light factors. The light factors are an important part of the analysis, and are parameterised as LF, the fraction of the total system light produced by one star for the wavelength or wavelength interval under consideration. The LFs for the binary components should sum to unity. The determination of the atmospheric parameters represents a difficult optimization problem, for which we use a genetic algorithm to minimise the χ^2 of the fit to the data (Holland 1975; Charbonneau 1995). Our implementation is called GENFITT (GENetic FITTING) and in approach is similar to that of Mokiem et al. (2005). Error estimates come from the covariance matrix, which is constructed using the Levenberg-Marquardt method.

In order to save computing time we pre-calculate grids of synthetic spectra. GENFITT linearly interpolates between these in T_{eff} and $\log g$, and then convolves them with a rotational profile using the ROTINF code of I. Hubeny¹. The LTE grid covers T_{eff} from 6000 to 15 000 K and the non-LTE grid covers T_{eff} from 14 000 to 35 000 K. Both grids contain $\log g$ values of 2.5 to 5.0 (cgs), and are stepped by 250 K in T_{eff} and 0.1 dex in $\log g$.

There can be strong degeneracies between the fitted atmospheric parameters, most notably T_{eff} and $\log g$ for Balmer line profiles. This degeneracy between T_{eff} and $\log g$ can be avoided by analysing dEBs, because their surface gravities can be known

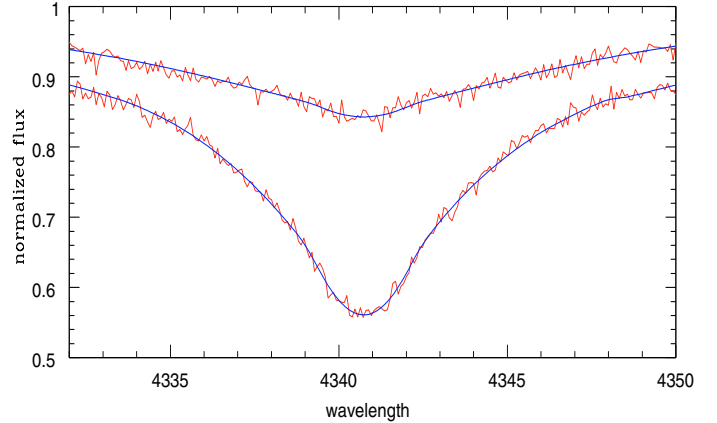


Fig. 1. Synthesized H γ line profiles (red lines) and the best-fitting synthetic spectra (blue lines) for the case with $S/N = 100$ in Table 1. The primary star is below the secondary star on the plot.

to within 0.01 dex from measurements of their masses and radii. In many cases the light ratio of the stars in a dEB can be obtained from the light curve analysis (e.g. Southworth et al. 2004b), so the LFs can then be fixed to known values².

In order to test the performance of GENFITT we synthesized disentangled spectra covering the H β , H γ and H δ lines, using representative atmospheric parameters and with Gaussian noise added to produce S/N ratios ranging from 25 to infinity. We then used GENFITT to fit all the atmospheric parameters to the synthesized spectra, with the only constraint that the two light factors sum to unity. The results are given in Table 1 and show that for high S/N ratios (≥ 100) the GENFITT results reproduce the input T_{eff} and $\log g$ values satisfactorily. For lower S/N ratios the inherent degeneracy of these parameters causes them both to be underestimated by our method. By contrast, the LFs are reproduced to well within the errorbars for all S/N ratios considered. An example fit is shown in Fig. 1

We obtained a second set of solutions in an “unconstrained mode” where the LFs were not required to sum to unity. The results were, as expected, similar to but slightly poorer than the “constrained mode”. Finally, a third set of solutions were made with fixed $\log g$ values (Table 2), as would often be the case when analysing a dEB. We find that the situation is similar to that for the first set of solutions. The main limitation on the quality of these results is the degeneracy between T_{eff} and $\log g$: we find a correlation coefficient of 0.98 between these parameters for both components. The correlation with the LFs is much weaker, which is why the LF values are reliable even for low- S/N spectra.

3. Application to V615 Persei

GENFITT has already been used for studying the dEBs V380 Cyg (Pavlovski et al. 2009) and V621 Per (Southworth et al. 2011, in prep.). In both of these cases spectra of a high S/N ratio were available and GENFITT returned excellent results. Here we challenge it with spectra of a much lower S/N ratio.

25 spectra of the dEB V615 Per were obtained by SMS04, with a reciprocal dispersion of $0.11 \text{ \AA pixel}^{-1}$, a resolution of $0.2 \text{ \AA}$ and an average S/N of ≈ 50 . They cover 4220–4500 Å so include H γ , a number of helium lines, and the Mg II 4481 Å doublet and are well distributed through one orbital cycle. SPD

² The light ratio can be poorly determined in some dEBs, and in this case obtaining a spectroscopic light ratio is an important part of modelling the light curves. For an example see Southworth et al. (2007).

¹ <http://nova.astro.umd.edu/index.html>

Table 1. Results of fitting synthesized disentangled spectra with GENFITT for different S/N values.

S/N	$T_{\text{eff}1}$ (K)	$\Delta T_{\text{eff}1}$	$\log g_1$	$\Delta \log g_1$	LF_1	ΔLF_1	$T_{\text{eff}2}$ (K)	$\Delta T_{\text{eff}2}$	$\log g_2$	$\Delta \log g_2$	LF_2	ΔLF_2	χ^2_ν
∞	17 991	−9	3.495	−0.005	0.775	−0.005	11 005	−5	3.995	−0.005	0.225	+0.005	−
	± 56		± 0.016		± 0.012		± 56		± 0.016		± 0.012		
500	17 896	−104	3.492	−0.008	0.771	−0.009	10 896	−103	3.992	−0.008	0.229	+0.009	1.054
	± 153		± 0.025		± 0.024		± 157		± 0.023		± 0.021		
200	17 835	−165	3.497	−0.003	0.771	−0.009	10 837	−164	3.987	−0.013	0.229	+0.009	1.032
	± 198		± 0.027		± 0.015		± 211		± 0.024		± 0.019		
100	17 753	−247	3.485	−0.015	0.769	−0.011	10 748	−252	3.986	−0.015	0.231	+0.011	0.841
	± 365		± 0.033		± 0.023		± 358		± 0.036		± 0.021		
50	17 041	−959	3.421	−0.079	0.761	−0.019	10 030	−970	3.929	−0.071	0.239	+0.019	0.774
	± 441		± 0.035		± 0.081		± 430		± 0.036		± 0.086		

Notes. The input atmospheric parameters are: $T_{\text{eff}1} = 18\,000$ K, $\log g_1 = 3.5$, $v \sin i_1 = 105$ km s $^{-1}$, $\text{LF}_1 = 0.78$, $T_{\text{eff}2} = 11\,000$ K, $\log g_2 = 4.0$, $v \sin i_2 = 95$ km s $^{-1}$, $\text{LF}_2 = 0.22$. The output parameters and uncertainties are given, along with the difference compared to the input parameters and the reduced χ^2 (χ^2_ν).

Table 2. Same as Table 1 but for fits with $\log g$ values for the two stars fixed.

S/N	$T_{\text{eff}1}$ (K)	$\Delta T_{\text{eff}1}$	LF_1	ΔLF_1	$T_{\text{eff}2}$ (K)	$\Delta T_{\text{eff}2}$	LF_2	ΔLF_2	χ^2_ν
∞	17 921	−79	0.776	+0.001	10 921	−79	0.224	−0.001	−
	± 109		± 0.010		± 115		± 0.014		
500	17 897	−103	0.771	−0.004	10 909	−91	0.229	+0.009	1.065
	± 160		± 0.022		± 166		± 0.021		
200	17 825	−175	0.769	−0.011	10 808	−192	0.231	+0.011	1.042
	± 236		± 0.018		± 236		± 0.018		
100	17 750	−250	0.767	−0.013	10 738	−262	0.233	+0.013	0.848
	± 352		± 0.027		± 343		± 0.025		
50	17 056	−944	0.756	−0.024	10 030	−970	0.244	+0.024	0.777
	± 416		± 0.082		± 422		± 0.085		

Table 3. Results from GENFITT analysis of the Hy line for V615 Per, with a comparison to the values found by SMS04.

Parameter	This work (solution A)		This work (solution B)		SMS04	
	Star A	Star B	Star A	Star B	Star A	Star B
T_{eff} (K)	14 710 $\pm$ 210	11 520 $\pm$ 290	14 920 $\pm$ 190	11 420 $\pm$ 250	15 000 $\pm$ 500	11 000 $\pm$ 500
$\log g$ [cgs]	4.302 $\pm$ 0.035	4.361 $\pm$ 0.030	4.328 fixed	4.381 fixed	4.328 $\pm$ 0.059	4.381 $\pm$ 0.050
LF	0.676 $\pm$ 0.006	0.324 $\pm$ 0.006	0.677 $\pm$ 0.004	0.323 $\pm$ 0.004	0.65 $\pm$ 0.03	0.35 $\pm$ 0.03

was performed in Fourier space using the FDBINARY³ code (Ilijić et al. 2004). The disentangled spectra have S/N values of about 160 for the primary star (star A) and 80 for the secondary (star B).

Since the spectra of V615 Per cover only a limited wavelength range, our estimate of the T_{eff} s of the component stars is restricted to the Hy line. Helium lines are also good T_{eff} indicators, but instead we will use these later to obtain the helium abundance of the binary. Because the available light curves of V615 Per are not definitive, the $\log g$ values for the two stars are known to modest accuracies of 0.059 and 0.050 dex. We therefore included them as GENFITT fitted parameters (solution A), along with T_{eff} . We also obtained a solution B for comparison, where the $\log g$ values were fixed to those found by SMS04. The LFs were constrained to sum to unity. For $v \sin i$ we adopted 28 ± 5 and 8 ± 5 km s $^{-1}$ plus an instrumental broadening of 16 km s $^{-1}$ (SMS04). The results are given in Table 3.

For the chemical abundance analysis we adopted the atmospheric parameters from solution A. Synthetic spectra were calculated using ATLAS9 model atmospheres (Kurucz 1979) and non-LTE theoretical line profiles from the DETAIL and SURFACE codes (Giddings 1981; Butler 1984). A canonical microturbulence velocity of 2 km s $^{-1}$ was adopted (Trundle et al. 2007) as

Table 4. Chemical abundances derived for the components of V615 Per.

Component	Species	Wavelength (Å)	Abundance
Star A	He I	4388	0.090 $\pm$ 0.005
Star A	He I	4471	0.092 $\pm$ 0.005
Star A	He I	mean	0.091 $\pm$ 0.007
Star A	Mg II	4481	7.26 $\pm$ 0.03
Star B	Mg II	4481	7.16 $\pm$ 0.06

we have too few spectral lines to fit for it. The helium abundance of star A was derived via χ^2 minimisation between the observed profiles of He I 4388 Å and 4471 Å, and profiles calculated for abundances in the range $\epsilon(\text{He}) = 0.06$ –0.15. The results for both lines are given in Table 4 and correspond to a mean helium abundance of $\epsilon(\text{He}) = 0.091 \pm 0.007$. This abundance is solar to within the uncertainty [$\epsilon_\odot(\text{He}) = 0.089$], so no deviations in helium abundance are detected for V615 Per A. The helium lines from star B are too weak to be useful (Fig. 2).

Magnesium is an excellent metallicity indicator for B-type stars as it does not participate in the CNO process so is unmodified by stellar evolution (Lyubimkov et al. 2005). These authors found a mean abundance of $\log \epsilon(\text{Mg}) = 7.59 \pm 0.15$ for nearby B stars with reliable microturbulence velocities, in excellent agreement with the solar value of $\log \epsilon_\odot(\text{Mg}) = 7.55 \pm 0.02$.

³ <http://sail.fer.zep/fdbinary/>

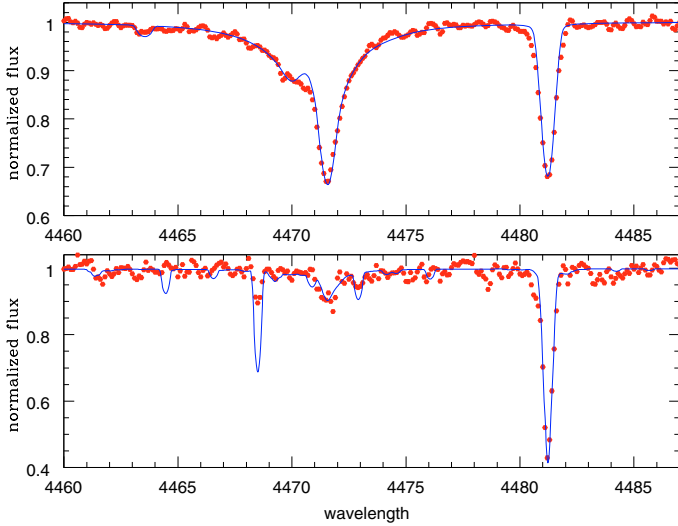


Fig. 2. Comparison between disentangled and renormalised spectra of star A (*upper panel*) and star B (*bottom panel*) with synthetic spectra calculated for the He and Mg abundances given in Table 4 (blue lines). Only the Mg II lines and the He I line for star A are fitted here. Other lines are shown but not used in our results.

Mg II 4481 Å is a prominent feature in the spectra of both components of V15 Per. Theoretical line profiles were calculated in non-LTE for star A and in LTE for star B and the abundances obtained by χ^2 minimisation (Fig. 2). We find a mean abundance of $\log \epsilon(\text{Mg}) = 7.21 \pm 0.07$. This is 0.46 dex lower than the mean value found by Lyubimkov et al. (2005) and 0.16 dex lower than that found by Daflon et al. (2003) for OB stars in the solar circle. The components of V615 Per have a subsolar Mg abundance and resemble halo B stars more than nearby examples (Daflon et al. 2004). This implies that the h Persei open cluster has a subsolar metal abundance.

4. Summary

Spectral disentangling is a method for obtaining the individual spectra of the components of a binary star system from composite spectra obtained at a range of orbital phases. A disadvantage of this method is that the continuum light ratios of the stars are not found, because this information is not present in the observed spectra without making assumptions about the spectral characteristics of the stars. We present the GENFITT program, which uses a genetic algorithm to fit synthetic spectra to the disentangled spectra of both components of a binary system simultaneously. It returns the best-fitting atmospheric parameters (T_{eff} and $\log g$) and the light contributions of the two stars. From tests with synthesized spectra we find that GENFITT performs extremely well in determining the light ratio. It also returns reliable T_{eff} and $\log g$ values in those cases where S/N of the input disentangled spectra is ≥ 100 , which is the usual situation for observational studies.

The light contributions of the two stars will normally sum to unity, which provides a useful constraint for GENFITT. Contaminating light from a third star can in principle be found, in cases when the light contributions of the two stars in the binary sum to less than unity. Once the light contributions of the stars have been found, their disentangled spectra can be renormalised to the correct continuum levels. The resulting spectra can then be analysed using standard methods for single stars. If the two stars are eclipsing, their surface gravity values may

be found to high precision and accuracy from analysis of the orbital velocity amplitudes found by spectral disentangling and light curves covering the eclipses.

As a demonstration of the method we applied GENFITT to spectra of the eclipsing system V615 Per, a member of the h Persei open cluster. The metal abundance of this cluster is controversial (see Sect. 1) but important in measuring its distance by the isochrone method. The spectra were disentangled and fed into GENFITT, and an abundance analysis was performed on the resulting renormalised spectra. The atmospheric parameters returned by GENFITT are in good agreement with previous work (SMS04) but are more precise. We find a normal solar helium abundance for V615 Per A (star B is cooler and has only weak helium lines). The magnesium abundances for both stars are lower than those found for nearby OB stars, indicating that h Persei has a subsolar metallicity. This is in agreement with the results of SMS04, based on the complementary method of comparing the masses and radii of the stars to the predictions of theoretical stellar evolutionary models.

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References

- Andersen, J. 1991, A&ARv, 3, 91
- Butler, K. 1984, PhD Thesis, University College London
- Capilla, G., & Fabregat, J. 2002, A&A, 394, 479
- Charbonneau, P. 1995, ApJS, 101, 309
- Currie, T., Hernandez, J., Irwin, J., et al. 2010, ApJS, 186, 191
- Daflon, S., Cunha, K., Smith, V. V., & Butler, K. 2003, A&A, 399, 525
- Daflon, S., Cunha, K., & Butler, K. 2004, ApJ, 606, 514
- Dufton, P. L., Brown, P. J. F., Fitzsimmons, A., & Lennon, D. J. 1990, A&A, 232, 431
- Giddings, J. 1981, PhD Thesis, University College London
- Hadrava, P. 1995, A&AS, 114, 393
- Hensberge, H., Pavlovski, K., & Verschuere, W. 2000, A&A, 358, 553
- Hensberge, H., Ilijic, S., & Torres, K. B. V. 2008, A&A, 482, 1031
- Holland, J. H. 1975, Adaptation in natural and artificial systems. an introductory analysis with applications to biology, control and artificial intelligence, ed. Holland, J. H.
- Ilijic, S., Hensberge, H., Pavlovski, K., & Freyhammer, L. M. 2004, in Spectroscopically and Spatially Resolving the Components of the Close Binary Stars, ed. R. W. Hilditch, H. Hensberge, & K. Pavlovski, ASP Conf. Ser., 318, 111
- Keller, S. C., Grebel, E. K., Miller, G. J., & Yoss, K. M. 2001, AJ, 122, 248
- Klochova, V. G., & Panchuk, V. E. 1987, Sov. Astron. Lett., 13, 23
- Krziesinski, J., Pigulski, A., & Kolaczowski, Z. 1999, A&A, 345, 505
- Kurucz, R. L. 1979, ApJS, 40, 1
- Lennon, D. J., Brown, P. J. F., & Dufton, P. L. 1988, A&A, 195, 208
- Lyubimkov, L. S., Rostopchin, S. I., Rachkovskaya, T. M., Poklad, D. B., & Lambert, D. L. 2005, MNRAS, 358, 193
- Marco, A., & Bernabeu, G. 2001, A&A, 372, 477
- Mokiem, M. R., de Koter, A., Puls, J., et al. 2005, A&A, 441, 711
- Nissen, P. E. 1976, A&A, 50, 343
- Pavlovski, K., & Hensberge, H. 2005, A&A, 439, 309
- Pavlovski, K., & Hensberge, H. 2009, ASP Conf., in press [arXiv:0909.3246]
- Pavlovski, K., & Southworth, J. 2009, MNRAS, 394, 1519
- Pavlovski, K., Tamajo, E., Koubicky, P., et al. 2009, MNRAS, 400, 791
- Simon, K. P., & Sturm, E. 1994, A&A, 281, 286
- Simon, K. P., Sturm, E., & Fiedler, A. 1994, A&A, 292, 507
- Slesnick, C. L., Hillenbrand, L. A., & Massey, P. 2002, ApJ, 576, 880
- Smart, S. J., & Rolleston, W. R. J. 1997, ApJ, 481, L47
- Southworth, J., & Clausen, J. V. 2007, A&A, 461, 1077
- Southworth, J., Maxted, P. F. L., & Smalley, B. 2004a, MNRAS, 349, 547
- Southworth, J., Maxted, P. F. L., & Smalley, B. 2004b, MNRAS, 351, 1277
- Southworth, J., Bruntt, H., & Buzasi, D. L. 2007, A&A, 467, 1215
- Torres, G., Andersen, J., & Gimenez, A. 2010, A&ARv, 18, 67
- Trundle, C., Dufton, P. L., Hunter, I., et al. 2007, A&A, 471, 625
- Vrancken, M., Lennon, D. J., Dufton, P. L., & Lambert, D. L. 2000, A&A, 358, 639
- Wolff, S. C., & Heasley, J. N. 1985, ApJ, 292, 589